Fast neutrino flavor conversions inside the neutrino sphere in the

core-collapse supernovae

MILAD DELFAN AZARI

Department of Pure and Applied Physics, Waseda University

Collaborators.

Shoichi Yamada, Taiki Morinaga, Hirotada Okawa, Wakana Iwakami @ Waseda University Hiroki Nagakura @ Princeton University, Shun Furusawa @ Tokyo Univ. of Science Akira Harada @ ICRR, University of Tokyo and Kohsuke Sumiyoshi @ National Institute of Technology

> M. Delfan Azari *et al.,* Phys. Rev. **D** 99, 103011 M. Delfan Azari *et al.,* (arXiv: 1910.06176)

Oct 21-24, 2019

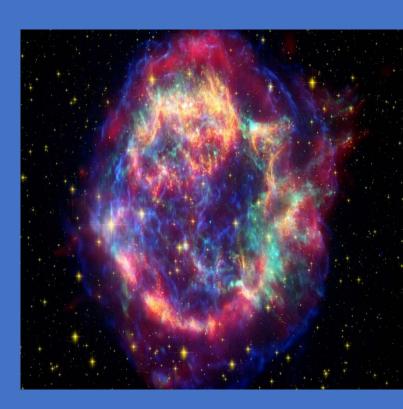
4M–COCOS, Fukuoka

Fate of a star depends on its Mass



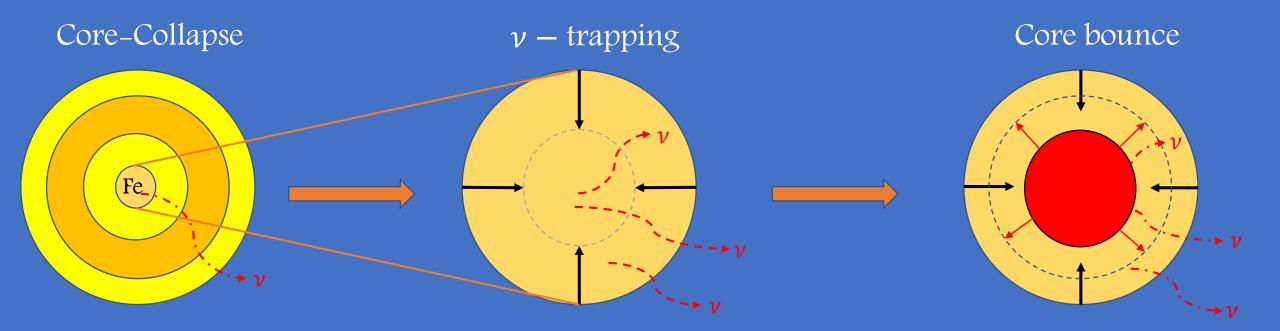
Core-Collapse

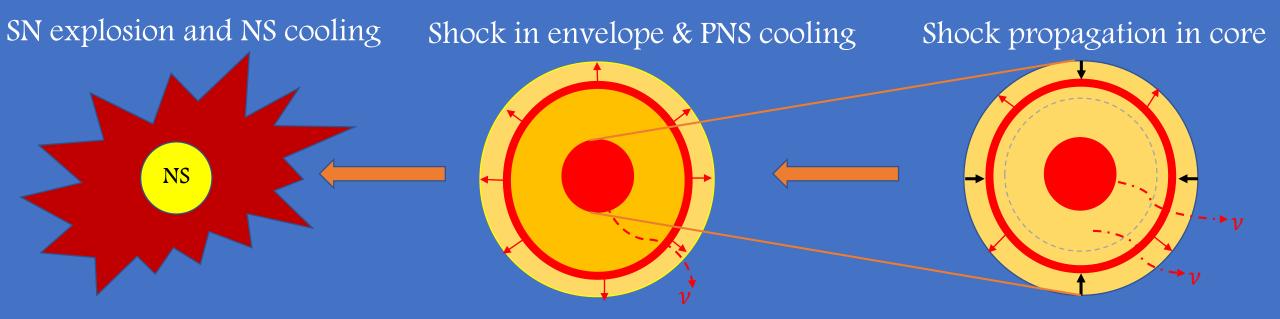
Supernovae



Cassiopeia A

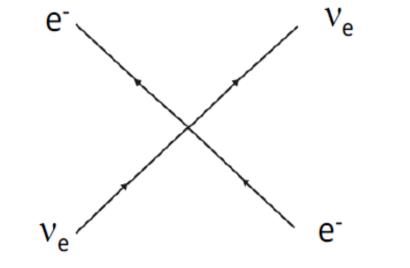


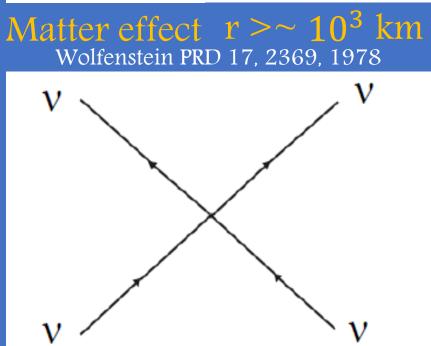




Why neutrinos?

- ✓ almost all of the binding energy of NS liberated in the gravitational collapse is emitted in the form of neutrinos and the kinetic energy of matter in the supernova explosion is just 1% of this energy.
- ✓ In the ν heating mechanism, a fraction of ν_e and $\bar{\nu}_e$ are re-absorbed by the matter between the shock front and the so-called gain radius and deposit their energy to push the stagnated shock again.





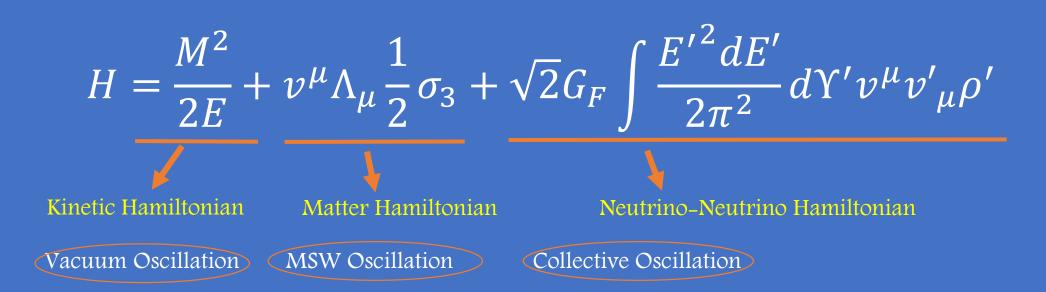
 $\nu - \nu$ self interaction $r \sim 10^2$ km Duan et al., PRD 74, 105014, 2006

Basic Equations and Formulae

Equation of Motion: $(\partial_t + v. \nabla_r)\rho = i[\rho, H]$

$$\rho = \frac{fv_e + fv_x}{2} + \frac{fv_e - fv_x}{2} \begin{bmatrix} S_{\mathbf{p}} & S_{\mathbf{p}} \\ S_{\mathbf{p}}^* & -S_{\mathbf{p}} \end{bmatrix}$$

 M^2 : Mass-squared matrix v^{μ} : (1,**v**) Λ^{μ} : $\sqrt{2}G_F(n_e - n_{e^+})u^{\mu}$ $d\Upsilon' = d\mathbf{v}'/4\pi$



Basic Equations and Formulae

Linearized Equation of Motion.

$$i(\partial t + v.\nabla_r)S_v = v^{\mu}(\Lambda_{\mu} + \Phi_{\mu})S_v - \int \frac{d\mathbf{v}'}{4\pi}v^{\mu}v'_{\mu}G_{v'}S_{v'}$$

$$G_{\nu} = \sqrt{2}G_F \int_0^\infty \frac{dEE^2}{2\pi^2} \left[f_{\nu_e}(E, \mathbf{v}) - f_{\overline{\nu}_e}(E, \mathbf{v}) \right]$$
$$\Phi^{\mu} \equiv \frac{d\mathbf{v}}{4\pi} G_{\nu} \nu^{\mu}$$

Assuming the solutions in the form of ;

 $S_{v}(t,r) = Q_{v}(\Omega, K)e^{-i(\Omega t - K.r)}$

$$\nu^{\mu}k_{\mu}Q_{\nu} = a^{\mu}$$

where ; $a^{\mu} \equiv -\int \frac{d\mathbf{v}'}{4\pi} v^{\mu} v'_{\mu} G_{v'} Q_{v'}$ $k^{\mu} = K^{\mu} - \Lambda^{\mu} - \Phi^{\mu}$ with $k^{\mu} = (\omega, \mathbf{k})$

 $\Pi^{\mu\nu}(\omega,\mathbf{k}) a_{\nu} = 0$

Polarization tensor $\Pi^{\mu\nu} = \eta^{\mu\nu} + \int \frac{dV}{4\pi} G_V \frac{v^{\mu}v^{\nu}}{\omega - V.k}$

 $D(\omega, \mathbf{k}) \equiv \det[\Pi] = 0$

Background Numerical Model

✓ Results of the realistic 2D simulations on the K-Supercomputer Nagakura et al., ApJ 854, 136 (2018)

✓ For non-rotating progenitor model of M_{star} =11.2 M_{\odot} Woosley et al., Reviews of Modern Physics 74, 1015 (2002)

✓ Boltzmann equation for neutrino transport being solved and special relativistic effect with a two energy grid technique. Nagakura et al., ApJS 214,16 (2014)

✓ Newtonian hydrodynamical equations & the Poisson equation for self-gravity were solved simultaneously.

Method

✓ Equations are written on spherical coordinates (r, θ) and three momentum space (*E*, θ_{ν} , ϕ_{ν})

✓ Computational domains :

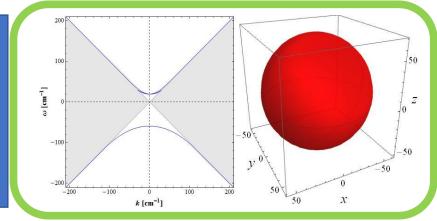
 $0 \le r \le 5000 \text{ km}$ $0 \le \theta \le \pi$ $0 \le E \le 300 \text{ MeV}$ $0 \le \theta_{\nu} \le \pi$ $0 \le \phi_{\nu} \le 2\pi$

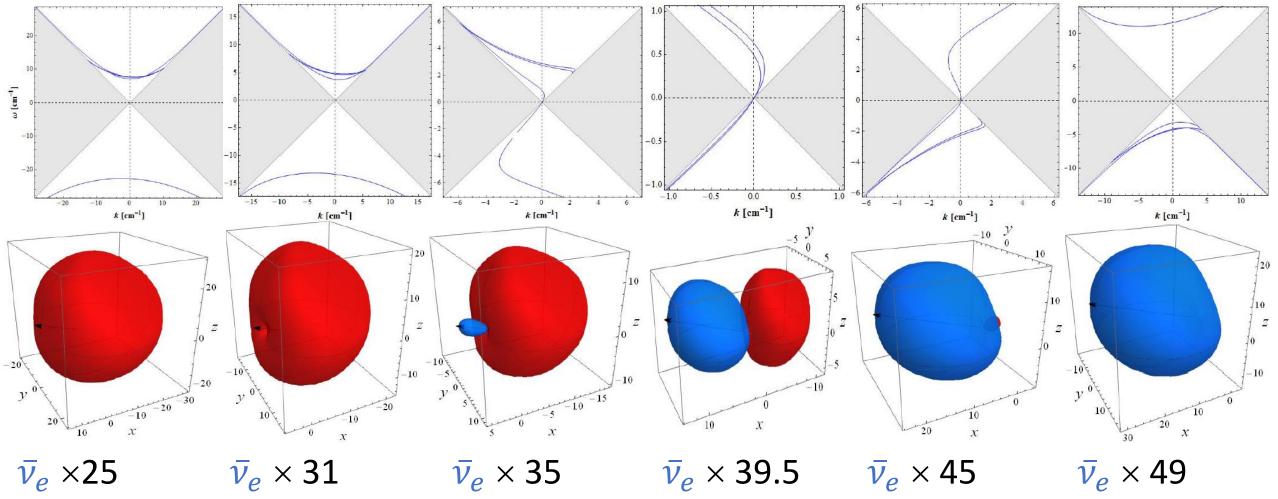
For 384 (r), 128 (θ), 20 (E), 10 (θ_{ν}) and 6 (ϕ_{ν}) mesh cells.

Results – Scaled data at t – 15 ms (r = 44.8 km, θ = 2.36 rad)

Crossing :

$$\int \frac{E^2 dE}{2\pi^2} \left(f_{\overline{\nu}_e}(E, \theta_{\nu}, \phi_{\nu}) - f_{\nu_e}(E, \theta_{\nu}, \phi_{\nu}) \right) \ge 0$$





New definitions

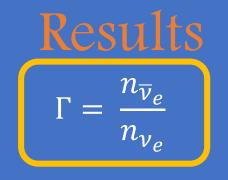
 Γ = The ratio of the anti-electron neutrino and the electron neutrino

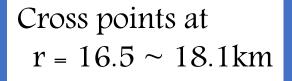
$$\Gamma = \frac{n_{\overline{\nu}_e}}{n_{\nu_e}}$$
$$n_{\nu} = \int \frac{d^3p}{(2\pi)^3} f_{\nu} \qquad \qquad d^3p = E^2 dE \sin\theta d\theta d\phi$$

Reminder :

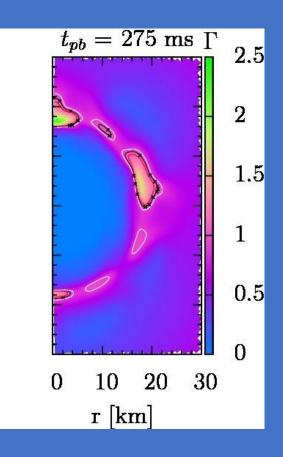
Crossing .

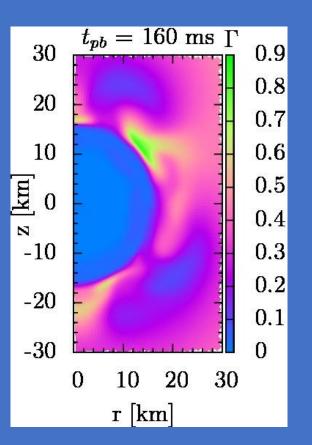
$$\int \frac{E^2 dE}{2\pi^2} \left(f_{\overline{\nu}_e}(E, \theta_{\nu}, \phi_{\nu}) - f_{\nu_e}(E, \theta_{\nu}, \phi_{\nu}) \right) \ge 0$$

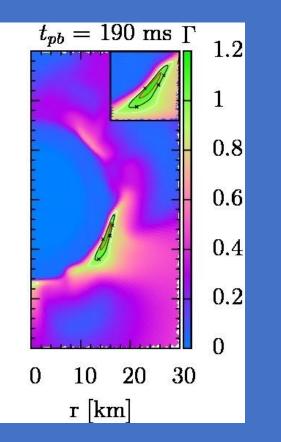




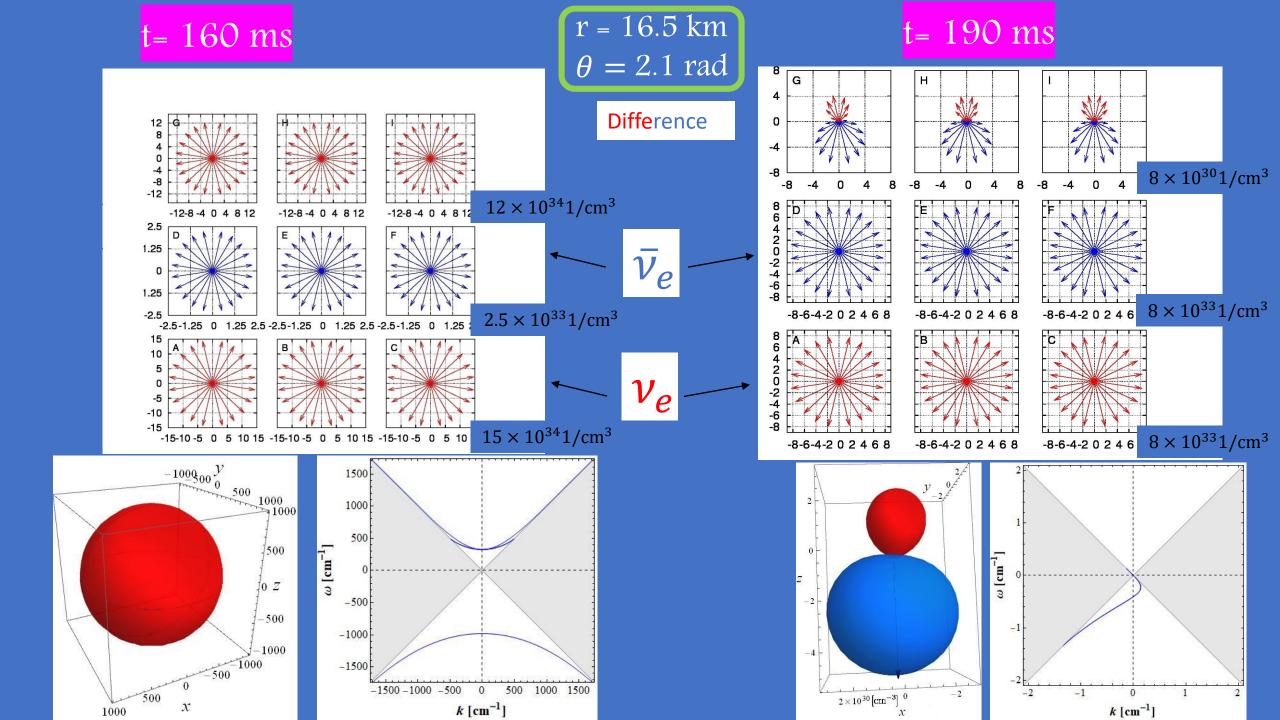
Cross points at $r = 16.1 \sim 21.4$ km

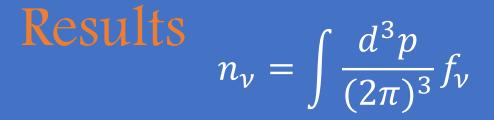


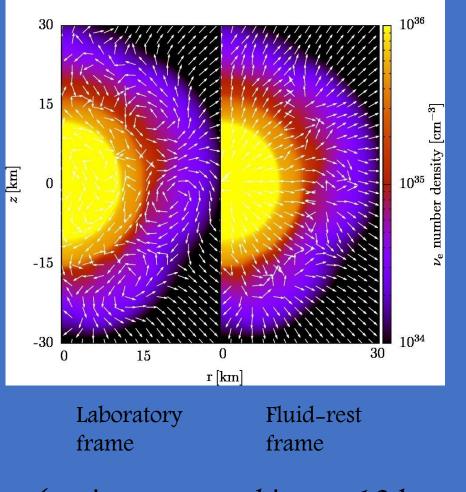




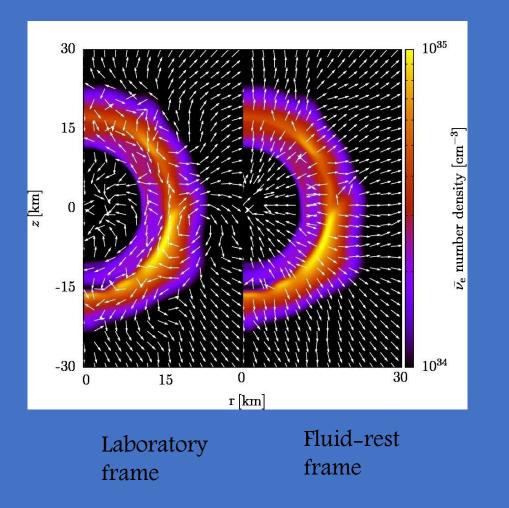
+ represents the crossing points





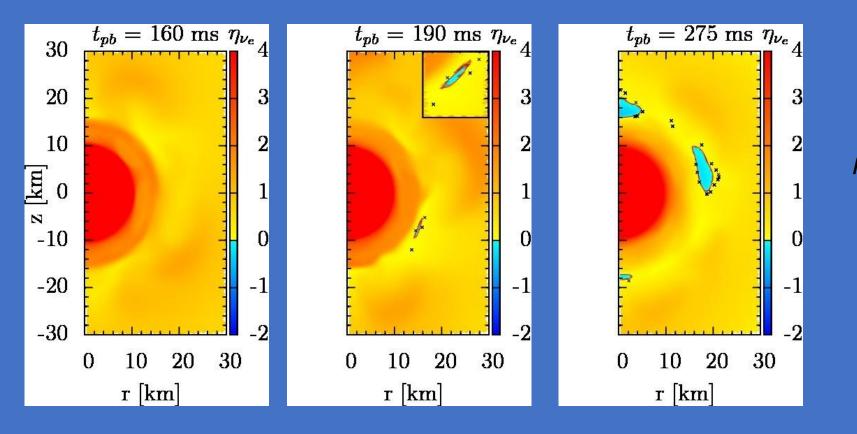


 $d^{3}p = E^{2}dE \sin\theta d\theta d\phi$



✓ \bar{v}_e is suppressed in r < 10 km. ✓ \bar{v}_e is more abundant at r ~ 15 km. Results

$$\eta_{\nu_e} = \frac{\mu_{\nu_e}}{T} = \frac{(\mu_e + \mu_p) - \mu_n}{T}$$



r = 16.5 km

$$\theta$$
 = 2.1 rad
 ρ = 2. 408 × 10¹³g/cm³
 Y_e = 0.13
 T = 20.4 MeV
 μ_e = 56.9 MeV
 μ_p = 848 MeV
 μ_p = 903 MeV

+ represents the crossing points

Summary and future works

- ✓ The possible oscillation conditions in different radii and time steps has been studied.
- ✓ A systematic survey for all radii based on the angular distributions difference between the electron-neutrinos and anti-electron type neutrinos.
- ✓ It is shown that the first crossing occurs at t ~ 190 ms and the crossings occur when electron-neutrinos and anti-electron neutrinos ratio (Γ) is equal to 1.
- ✓ As recently some data for the 3D simulations from my colleagues is available, the similar investigations might be a point of interest.
- ✓ This method might be applicable to other astrophysical objects such as NS-mergers disks, etc.
 We plan to investigate the conversion possibilities in these objects in near future.