

The formation of Magnetars: from protoneutron star dynamos to core-collapse supernovae central engines

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Outline

Introduction

- Magnetars and supernovae
- Modeling

Magnetorotational instability

- Characteristics of the MRI
- Numerical models

Convective dynamo

- Convection in proto-neutron stars
- Numerical results

Magnetorotational explosions

- Non-dipolar magnetic fields
- Numerical models

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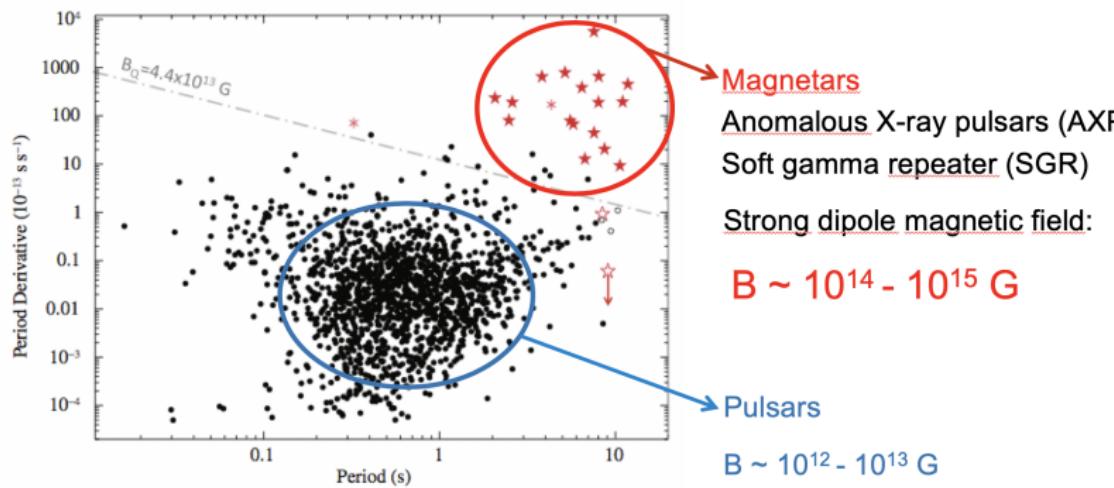
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Magnetars: the most intense known magnetic fields



Which supernovae are associated with magnetars?

Outstanding explosions and magnetic fields

Explosion kinetic energy

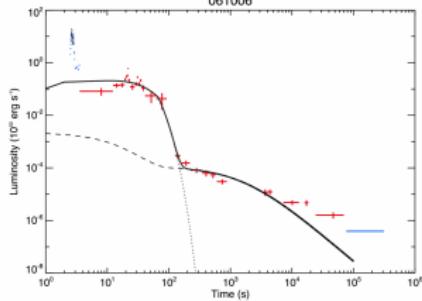
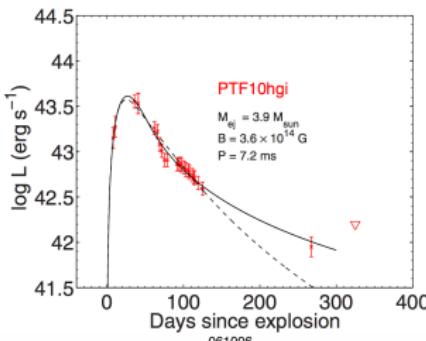
- ▶ Typical supernova: 10^{51} erg
- ▶ Rare hypernovae and GRB: 10^{52} erg

Total luminosity

- ▶ Typical supernova: 10^{49} erg
- ▶ Superluminous SN: 10^{51} erg

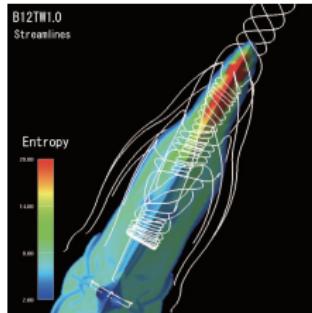
Lightcurves and X-ray plateaus

- ▶ Strong dipolar magnetic field:
 $B \sim 10^{14} - 10^{15}$ G
- ▶ Fast rotation: $P \sim 1 - 10$ ms
- ▶ Kasen and Bildsten (2010); Dessart et al. (2012); Nicholl et al. (2013);
 Inserra et al. (2013); Zhang and Mészáros (2001); Metzger et al. (2008);
 Lü et al. (2015); Gao et al. (2016)

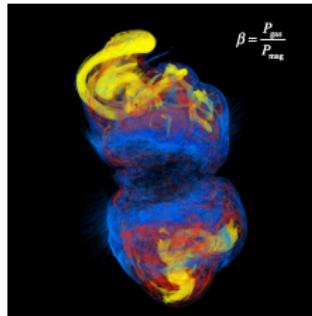


Impact of strong magnetic fields

- ▶ Strong magnetic fields ($B \sim 10^{15}$ G) and fast rotation ($P \sim$ few ms)
- ▶ Powerful jet-driven explosions (Shibata et al., 2006; Burrows et al., 2007; Dessart et al., 2008; Takiwaki et al., 2009; Kuroda and Ueda, 2010; Winteler et al., 2012; Obergaulinger and Aloy, 2017)
- ▶ In 3D the jet might be destroyed by **kink instability** (Mösta et al., 2014a)
- ▶ Big uncertainties on the **initial magnetic field**



Takiwaki et al. 2009



Moësta et al. 2014

Open question: origin of the magnetic field

- ▶ Compression of progenitor field in CCSN: $10^{12} - 10^{13}$ G (?)
- ▶ Magnetic field of NS before merger: $10^8 - 10^{12}$ G (but see Schneider et al. 2019)
- ▶ Magnetar: 10^{15} G

What is the amplification mechanism?

Possible dynamics within the PNS:

Magnetorotational Instability

Convective dynamo

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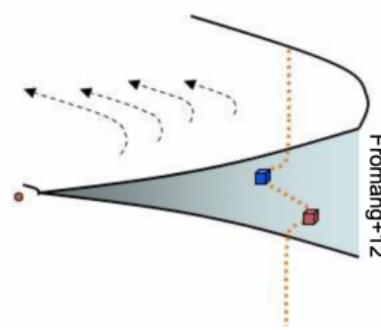
Numerical models



The magnetorotational instability (MRI)

In ideal MHD (no resistivity or viscosity)

- ▶ Instability onset: $\frac{d\Omega}{dr} < 0$
- ▶ Growth rate: $\sigma = \frac{q}{2}\Omega$ with $\Omega \propto r^{-q}$
- ▶ **Fast growth for fast rotation**
- ▶ Wavelength: $\lambda \propto \frac{B}{\sqrt{\rho\Omega}}$
- ▶ **Short wavelength for weak magnetic fields**



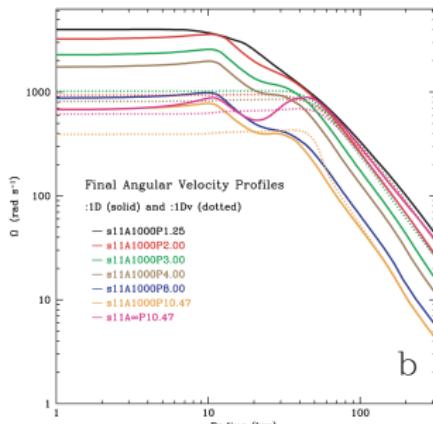
Proto-neutron stars vs. accretion disks

MRI unstable profile at $r > 10$ km

(Akiyama et al., 2003; Obergaulinger et al., 2009)

Conditions specific to PNS:

- ▶ neutrinos
- ▶ buoyancy
- ▶ spherical geometry



Ott et al. (2006)

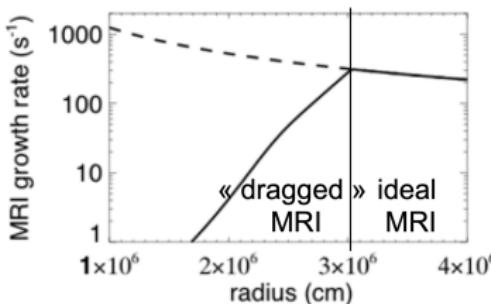
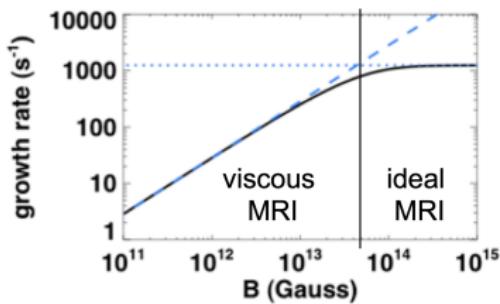
The impact of neutrinos on the MRI (Guilet et al., 2015, 2017)

Viscous regime

- ▶ $\lambda_{\text{MRI}} > l_\nu$
- ▶ Slow growth for weak fields

ν -drag regime

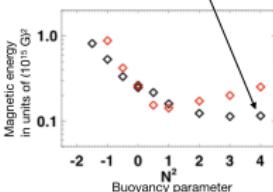
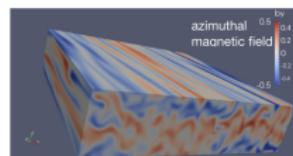
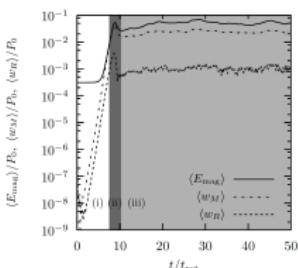
- ▶ $\lambda_{\text{MRI}} < l_\nu$
- ▶ Fast growth near the surface



Numerical simulations: local models

- ▶ Shearing box simulations
- ▶ Differential rotation with shearing boundary conditions
- ▶ Entropy/composition gradients
- ▶ Efficient amplification of the magnetic field, but no large-scale component
- ▶ Obergaulinger et al. (2009); Masada et al. (2012); Guilet et al. (2015); Rembiasz et al. (2016a,b)

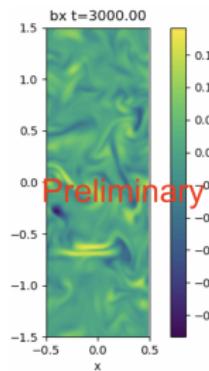
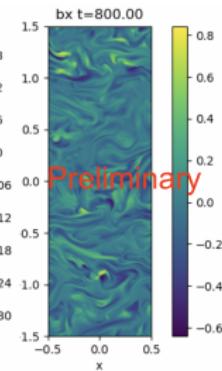
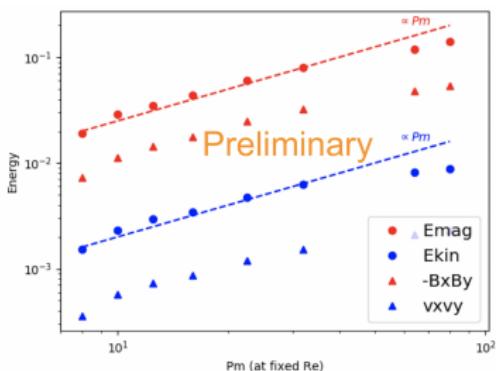
Masada et al. 2012



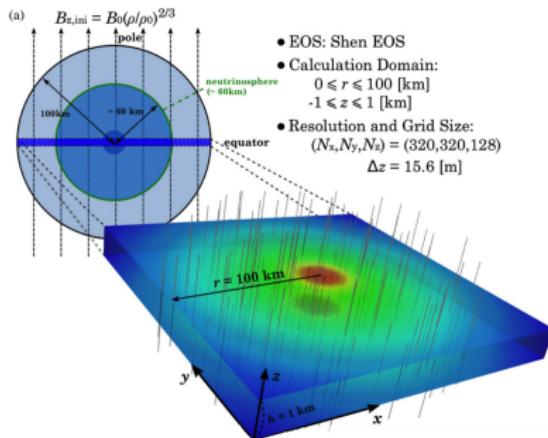
Guilet et al. 2015

Dependence on diffusion process

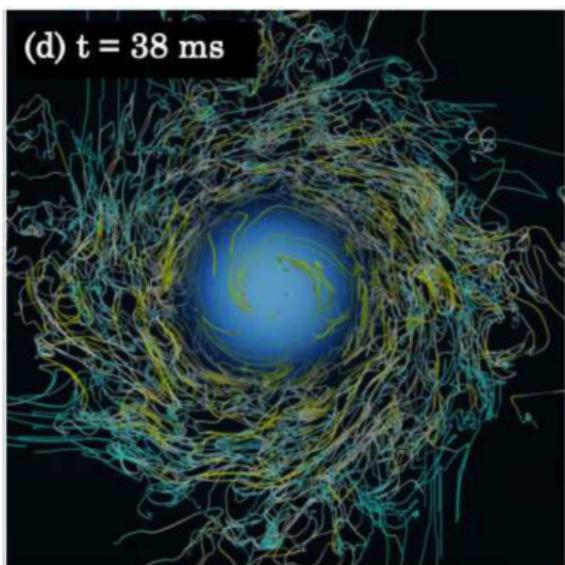
- ▶ Magnetic Prandtl number: $Pm = \nu/\eta$
- ▶ Proto-neutron star: $Pm \sim 10^{13}$
- ▶ Numerical simulations: $Pm \sim 1$
- ▶ Is there a **plateau at large Pm ?**

 $Pm=12.5$  $Pm=80$ 

Global models in the equatorial plane (Masada et al., 2015)

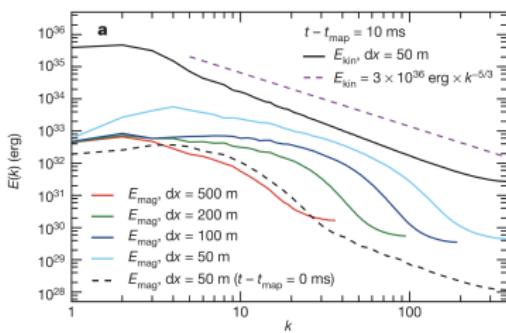
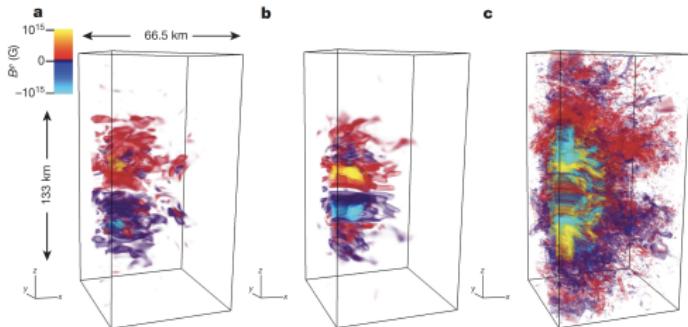


Magnetic field lines



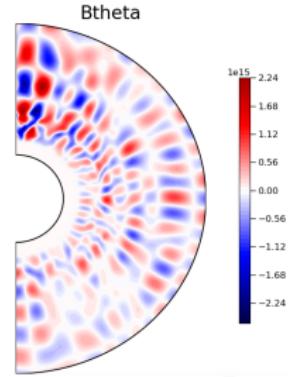
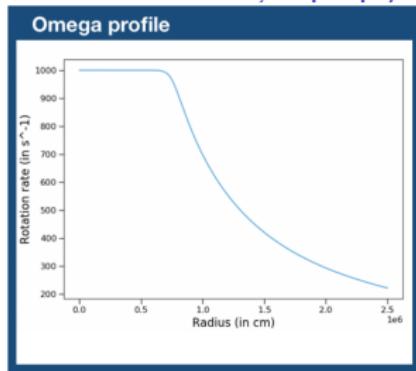
First attempt at a global model (Mösta et al., 2015)

- ▶ Global high-resolution (up to 50 m)
- ▶ Generation of large scale components
- ▶ Initial field: **strong dipole**

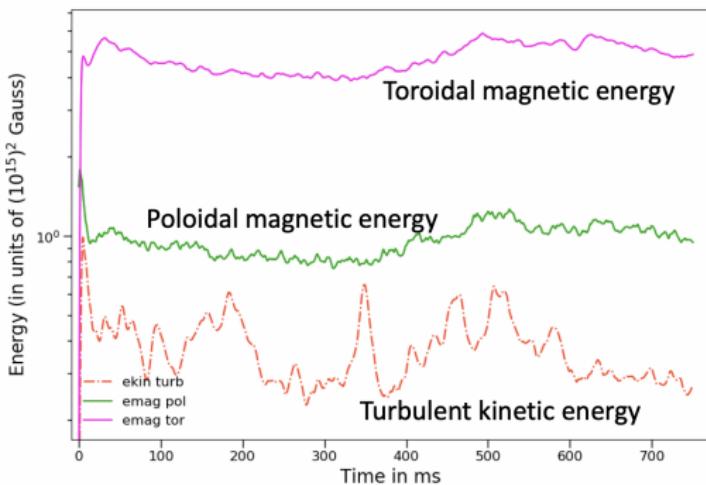
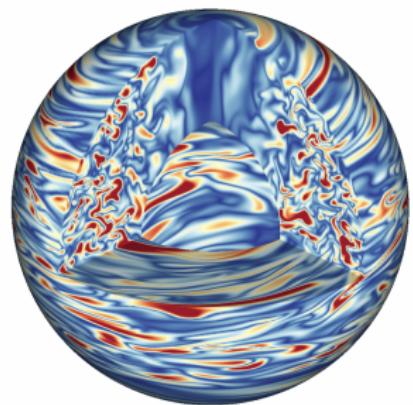


A simple setup for global MRI (Reboul-Salze et al., in prep.)

- ▶ Incompressible fluid
- ▶ Resolution: $256 \times 512 \times 1024$
- ▶ Typical parameters:
 - ▶ $Pm = 16$ and $Re = 5000$
 - ▶ $\Omega = 10^3 \text{ s}^{-1}$
 - ▶ $\nu = 8 \cdot 10^{11} \text{ cm}^2 \text{s}^{-1}$
 - ▶ $\eta = 5 \cdot 10^{10} \text{ cm}^2 \text{s}^{-1}$
 - ▶ $r = 25 \text{ km}$
 - ▶ $B_{\text{mean}} = 8.97 \cdot 10^{14} \text{ G}$



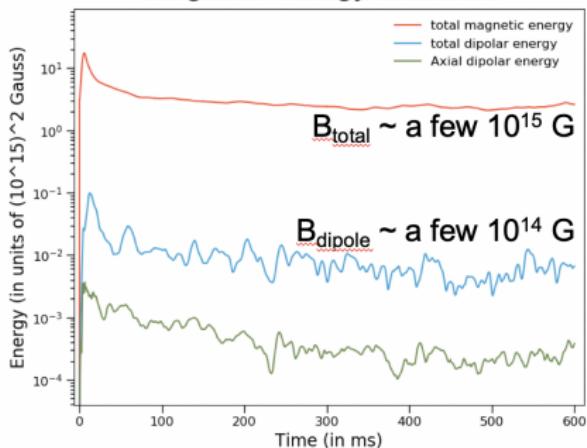
Geometry of the magnetic field



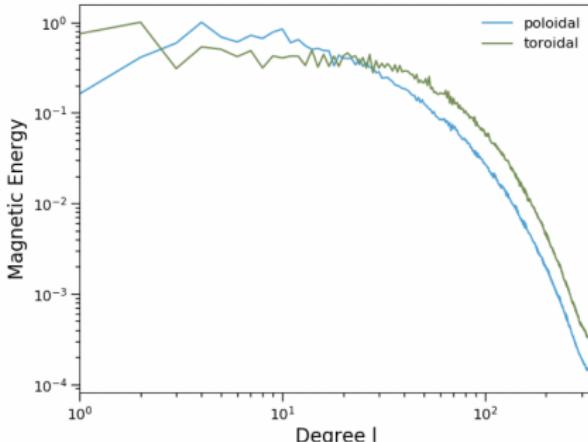
A. Reboul-Salze, J. Guilet, R. Raynaud & M. Bugli 2019, in prep.

A subdominant but significant dipole

Magnetic energy evolution



Magnetic energy spectrum



MRI saturated state independent of the initial magnetic field
Non-aligned dipole as strong as a (weak) magnetar

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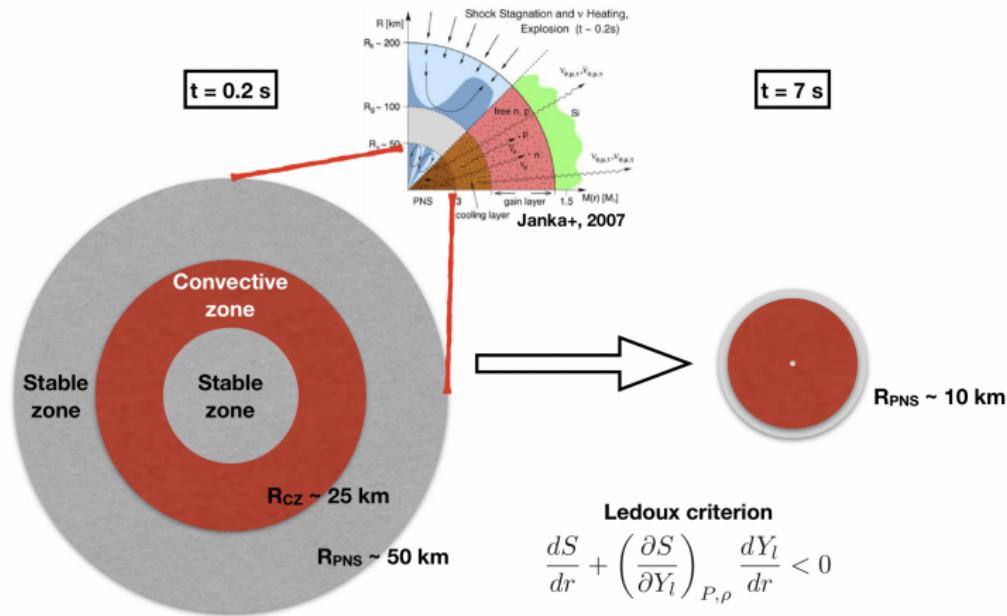
Numerical models





Convection in proto-neutron stars

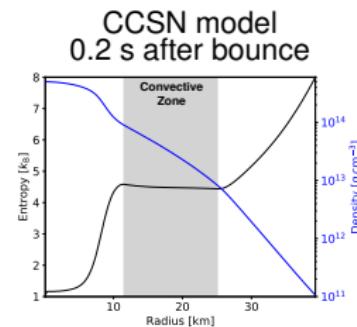
Convective zone in a proto-neutron star





Interior model and assumptions (Raynaud et al., submitted)

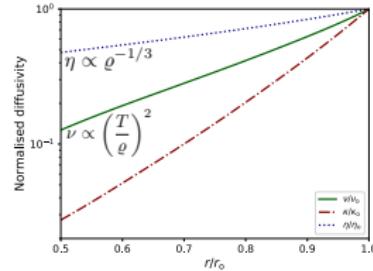
- ▶ PNS structure from 1D model by the Garching group
- ▶ Equation of state: LS220
- ▶ Progenitor: $27 M_{\odot}$
- ▶ Only convective zone: 12-25 km
- ▶ **Anelastic approximation** (soundproof)
- ▶ **Viscosity** and **thermal diffusion** by neutrinos
- ▶ **Resistivity** by electrons



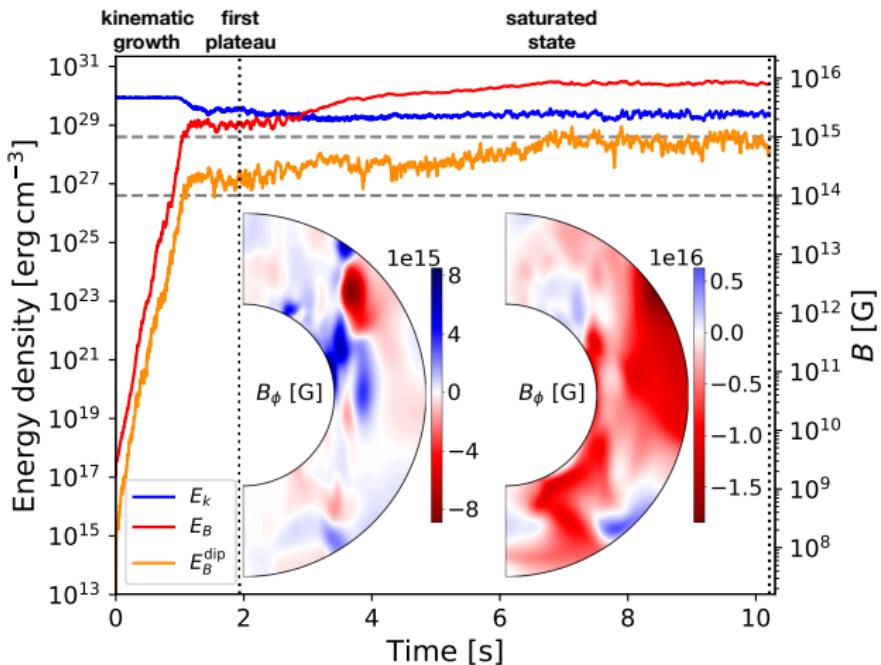
PNS

- ▶ $\nu_0 \sim 10^{10} \text{ cm}^2 \text{s}^{-1}$
- ▶ $\kappa_0 \sim 10^{12} \text{ cm}^2 \text{s}^{-1}$
- ▶ $\eta_0 \sim 10^{-5} \text{ cm}^2 \text{s}^{-1}$
- ▶ $\nu_0 \sim 10^{12} \text{ cm}^2 \text{s}^{-1}$
- ▶ $\kappa_0 \sim 10^{13} \text{ cm}^2 \text{s}^{-1}$
- ▶ $\eta_0 \sim 10^{12} \text{ cm}^2 \text{s}^{-1}$

Simulation



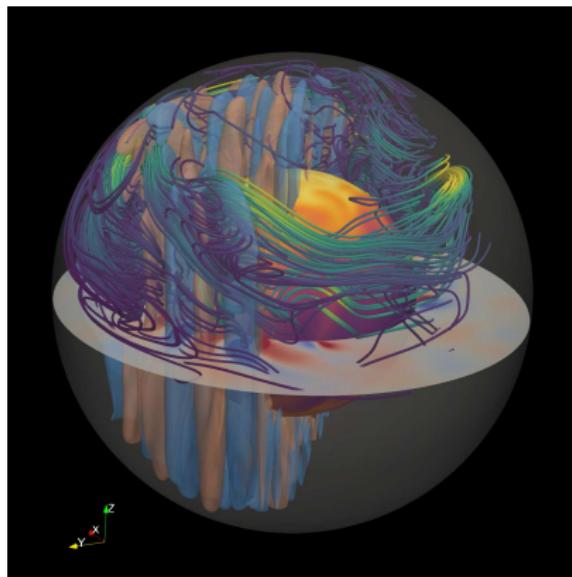
Typical time serie for fast rotating model ($P_{\text{rot}} = 2\text{ms}$)



Flow and magnetic field structure

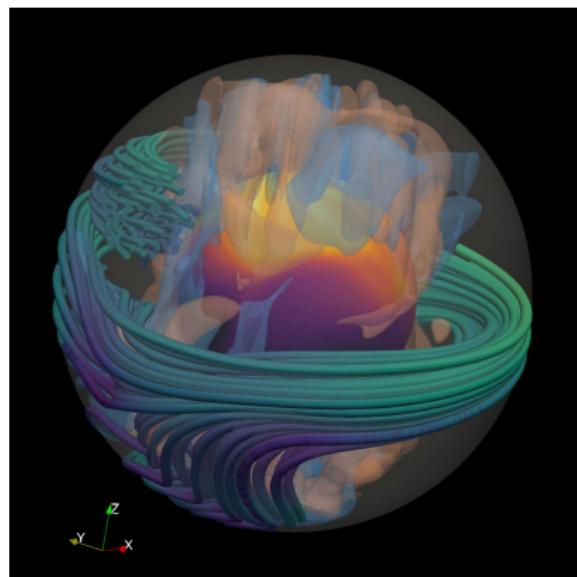
first plateau

$t \sim 2$ s



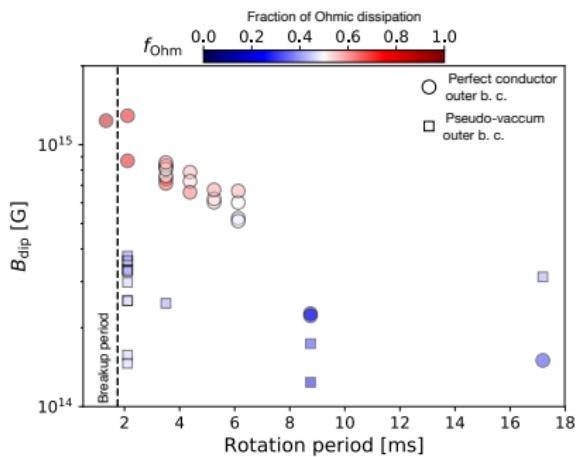
saturated state

$t \geq 6$ s

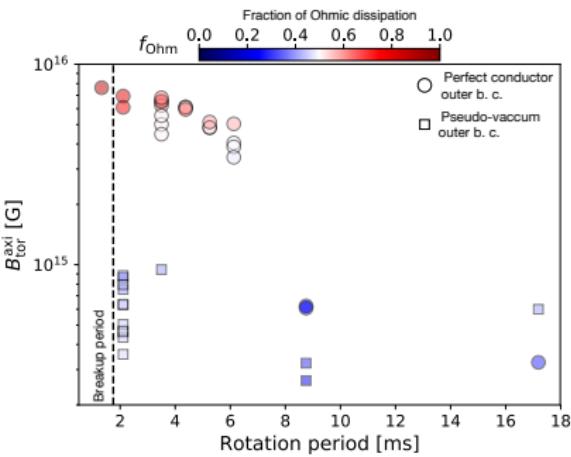


Magnetic field strength

B_{dip} up to $\sim 10^{15}$ G

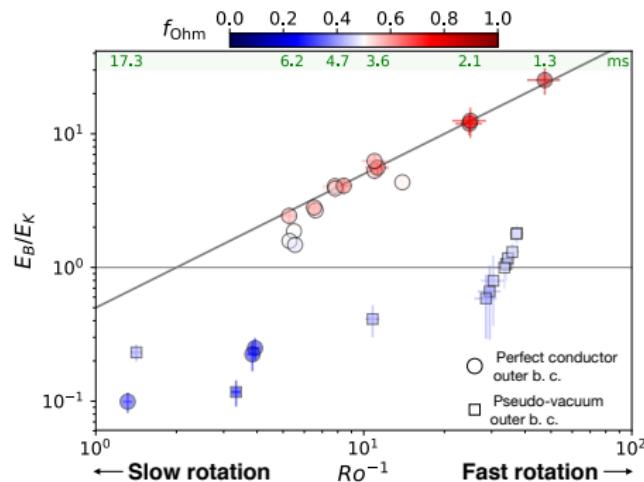


B_{tor} up to $\sim 10^{16}$ G



Values compatible with magnetar model for GRB

Understanding the dynamo saturation



Magnetostrophic
balance:
Lorentz-Coriolis

$$B^2 \propto \Omega \mu_0 U d$$



$$\frac{E_B}{E_{kin}} \propto Ro^{-1} = \frac{\Omega d}{U}$$

Delayed strong dynamo action?

- ▶ The PNS shrinks and cools down: higher spin, slower convective motions
- ▶ Rossby number Ro decreases: possible **transition to strong field regime**

Fast rotation

- ▶ $P < 2.5$ ms
- ▶ **Prompt strong dynamo**
- ▶ Hypernovae? GRBs?

Intermediate rotation

- ▶ $2.5 \text{ ms} < P < 10 \text{ ms}$
- ▶ **Delayed strong dynamo**
- ▶ SLSN? Magnetars in normal SNe?

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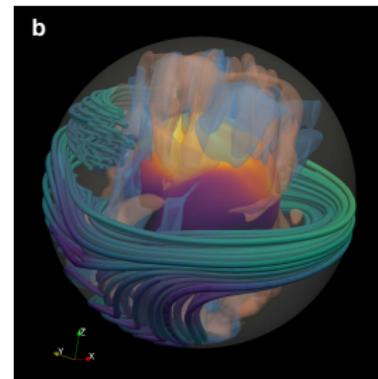
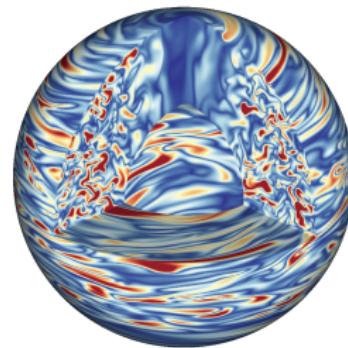
Non-dipolar magnetic fields

Numerical models



Connecting small and large scales

- ▶ Fields amplified to magnetar-like strength ($10^{14} - 10^{15}$ G).
- ▶ Topology: **non-dipolar** fields produced by dynamos.
- ▶ Time evolution: **delay** between the bounce and the rise of the magnetic field.
- ▶ At what stages can the amplified field affect the **shock revival**?



Initial magnetic field: pure dipole?

- ▶ Poor constraints from both observations and evolutionary models on the initial field.
- ▶ Quasi-uniform field up to $r_0 \sim 10^3$ km, then magnetic dipole (Suwa et al., 2007):

$$A_\phi = \frac{B_0}{2} \frac{r_0^3}{r^3 + r_0^3} r \sin \theta$$

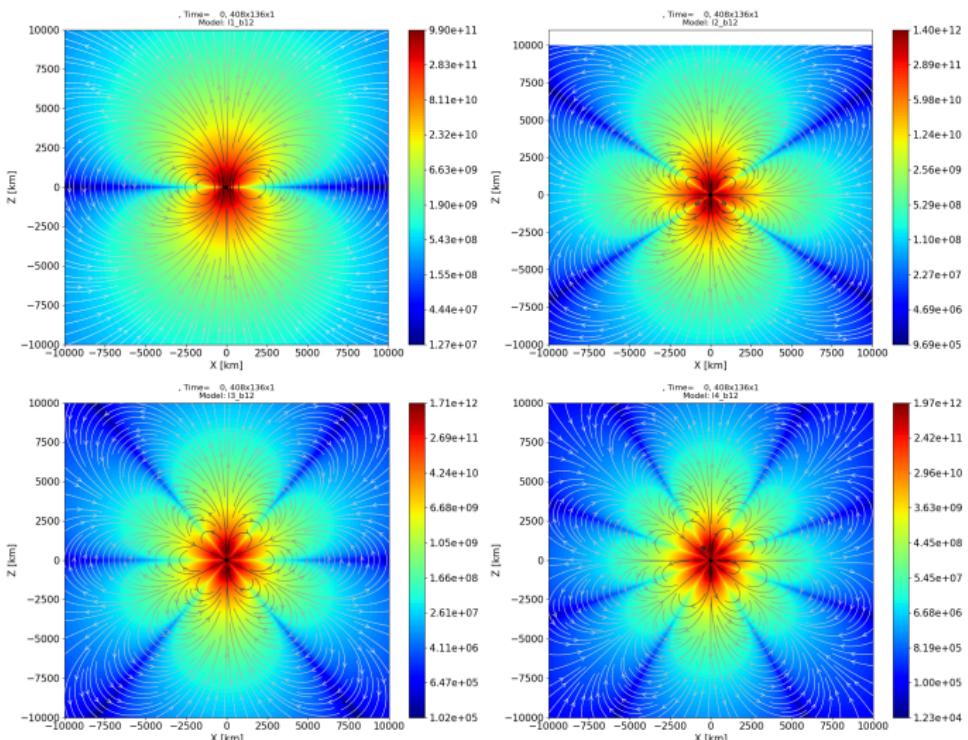
- ▶ Progenitor *original* field from 1D stellar evolution (Woosley and Heger, 2006)
- ▶ Superimposed **toroidal field** (Obergaulinger and Aloy, 2017):

$$B_\phi = B_0 \frac{r_0^3}{r^3 + r_0^3} r \cos \theta$$



Numerical models

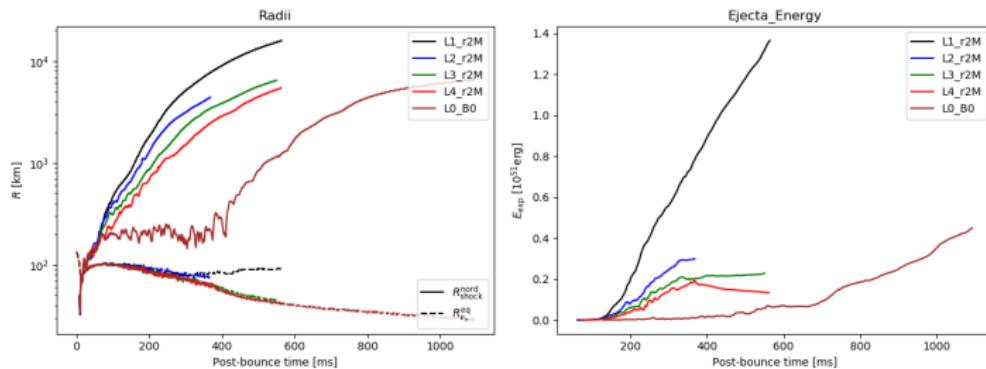
Comparison: different multipoles (Bugli et al., submitted)



Onset of the explosion ($r_0 = 2.9 \times 10^8$ cm, $B_0 = 8 \times 10^{10}$ G)

Shock radii - Explosion energies ($r_0 = 2.9 \times 10^8$ cm, $B_0 = 8 \times 10^{10}$ G)

- ▶ Onset of explosion at the same time
- ▶ **Slower expansion** for higher multipoles
- ▶ **Less energetic explosion** and **shallower increase of energy** for higher multipoles

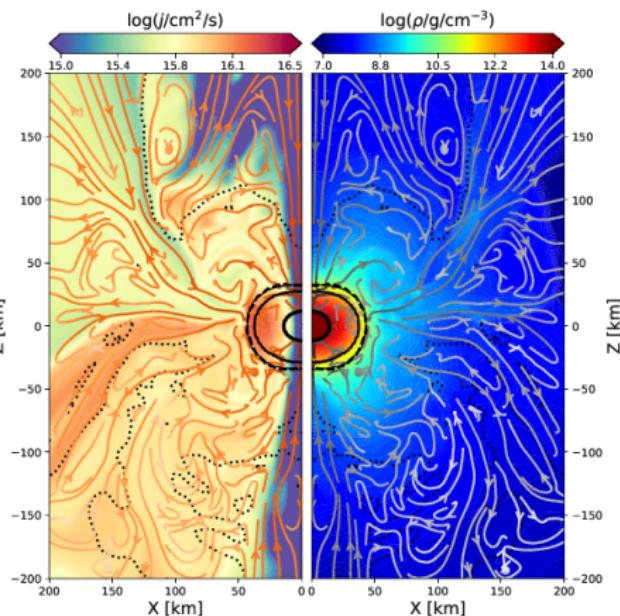
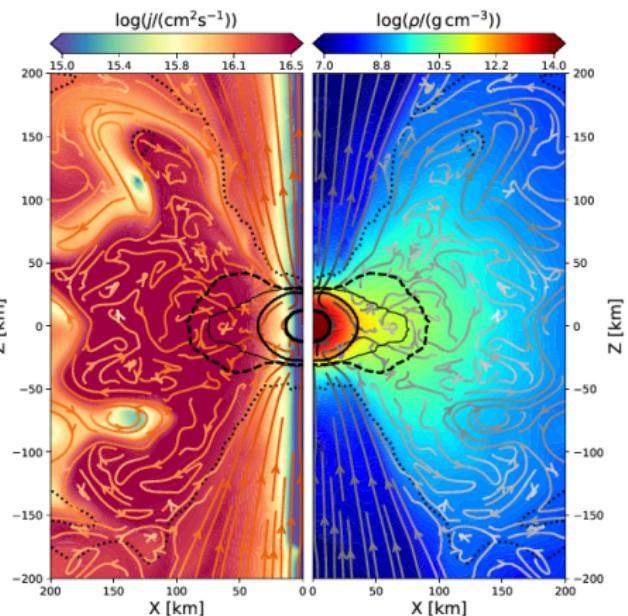


Magnetic/thermal pressure ($r_0 = 2.9 \times 10^8$ cm, $B_0 = 8 \times 10^{10}$ G)



Numerical models

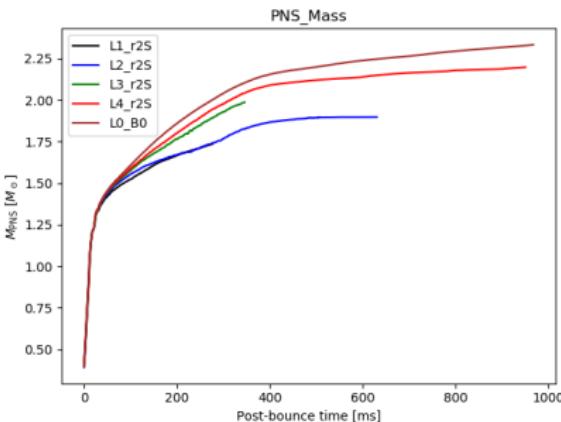
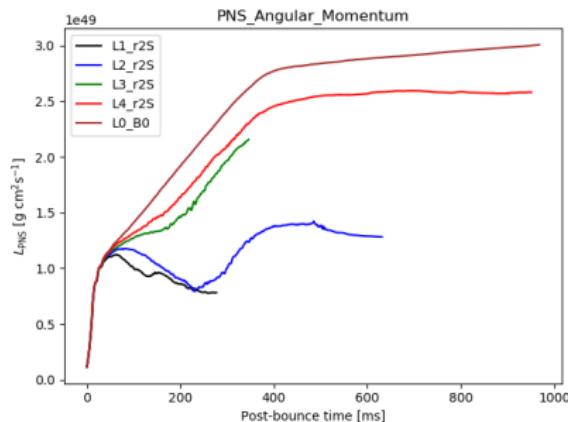
$I = 1$ vs $I = 4$ ($r_0 = 2.9 \times 10^8$ cm, $B_0 = 8 \times 10^{10}$ G)



PNS mass and spin ($r_0 = 2.9 \times 10^8$ cm, $B_0 = 2.6 \times 10^{11}$ G)

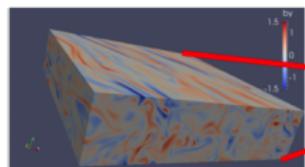
Higher multipole order \Rightarrow closer to hydrodynamic limit:

- ▶ faster increase of total angular momentum (weaker magnetic braking)
- ▶ faster increase of total mass



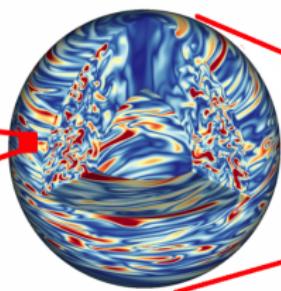
Still a long way to connect small and large scales!

Local Models (MRI)
 $L \sim 1-5$ km



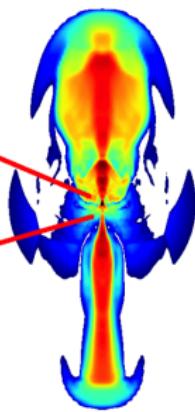
Guilet & Müller 2015
Guilet et al. 2015

Global Models
(Convection and MRI)
 $L \sim 10-50$ km



Raynaud et al. (submitted)
Reboul-Salze et al. (in prep.)

Magneto-rotational
supernova explosion
 $L \sim 100-10000$ km



Bugli et al. (submitted)

Conclusions

Magnetorotational instability

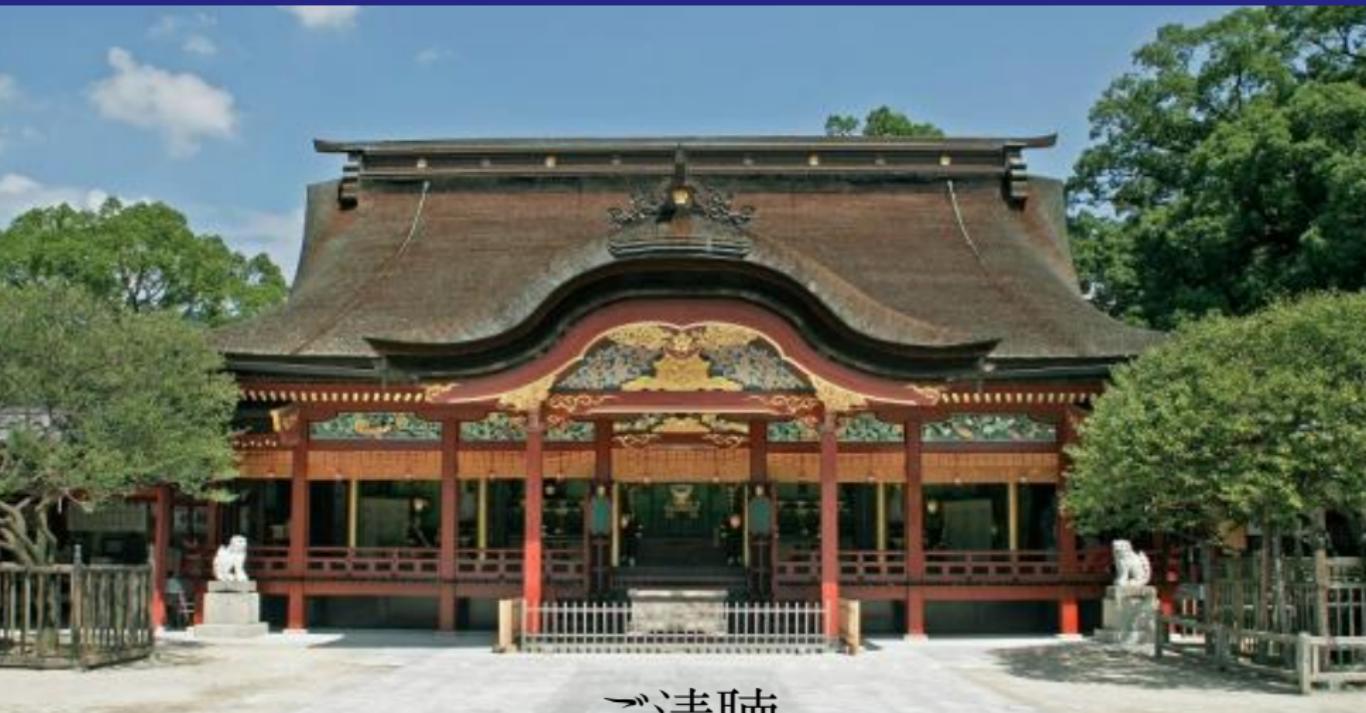
- ▶ Turbulent magnetic field with sub-dominant (but magnetar-like) dipole.
- ▶ Non-aligned dipole robustly generated from small scales.
- ▶ Buoyancy? Interaction with the convective dynamo?

Convective dynamos

- ▶ Amplification of seed fields to magnetar values.
- ▶ Clear dependence on the PNS rotation, possible transition in time to strong field regime.
- ▶ More dynamic evolution? Interaction with MRI unstable layer?

MHD-driven explosions

- ▶ Different topologies affect explosion dynamics and PNS formation.
- ▶ Lower multipoles produce more energetic explosions, slower and less massive PNS.
- ▶ 2D vs. 3D? Dynamo subgrid modeling?



ご清聴
ありがとうございました!



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