Well-posedness of renormalized solutions for a stochastic *p*-Laplace equation with *L*¹-initial data

Aleksandra Zimmermann University of Duisburg-Essen and University of Rostock (joint work with Niklas Sapountzoglou)

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 $T > 0, D \subset \mathbb{R}^d$ is a bounded Lipschitz domain, $Q_T := (0, T) \times D$. $du - \operatorname{div}(|\nabla u|^{p-2}\nabla u) dt = \Phi d\beta_t \quad \text{in } \Omega \times Q_T$ $u = 0 \quad \text{on } \Omega \times (0, T) \times \partial D$ $u(0, \cdot) = u_0$

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- $u_0 \in L^1(\Omega \times D)$ is \mathcal{F}_0 -measurable.
- ► The right-hand side is an integral in the sense of Itô with respect to $(\beta_t)_{t\geq 0}$, $\Phi \in L^2(\Omega \times Q_T)$ progressively measurable.

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- Randomness can be introduced as random external force by adding a stochastic integral on the right-hand side of the equation and by considering random initial values.
- The initial values may have poor regularity with respect to both variables.

Existence and uniqueness for L^2 - initial data

A unique strong solution can be obtained using classical monotonicity methods for SPDEs [Pardoux;1975], [Krylov, Rozovskii; 1983], [Liu, Röckner; 2015],...

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Let $u_0 \in L^2(\Omega \times D)$ be \mathcal{F}_0 -measurable. There exists a unique, \mathcal{F}_t -adapted, square-integrable stochastic process $u : \Omega \times [0, T] \rightarrow L^2(D)$ with a.s. continuous paths such that $u(0, \cdot) = u_0, u \in L^p(\Omega; L^p(0, T; V))$ and

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in $L^2(D)$ for all $t \in [0, T]$, a.s. in Ω .

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For p ≥
$$\frac{2d}{d+2}$$
, V = W₀^{1,p}(D).
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The Itô formula yields the energy estimate

$$\begin{split} &\frac{1}{2} \|u(t)\|_{L^2(D)}^2 - \frac{1}{2} \|u_0\|_{L^2(D)}^2 + \int_0^t \|\nabla u(s)\|_{L^p(D)}^p \, ds \\ &= \frac{1}{2} \int_0^t \|\Phi(s)\|_{L^2(D)}^2 \, ds + \int_0^t \int_D u(s) \Phi(s) \, dx \, d\beta_s. \end{split}$$

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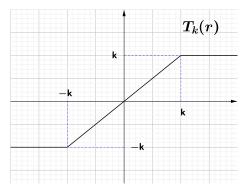
For *F*₀-measurable u₀ ∈ L²(Ω × D), the Itô formula yields the energy estimate

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- Renormalization: Generic concept of strategies to get rid of infinities.
- Renormalized solutions: Introduced by DiPerna and Lions 1989 in the study of the Boltzmann equation.
- ► Main idea: Nonlinear change of unknown v = S(u), where S is chosen in order to remove infinite quantities.

Renormalization and truncation

For $S \in C^2(\mathbb{R})$ such that $supp(S') \subset [-M, M]$ for some M > 0, S is constant outside [-M, M]. Thus, $S(u)(t) = S(T_k(u))(t)$ for all $k \geq M$.



For $u \in W_0^{1,p}(D)$, from the chain rule for Sobolev functions it follows that

$$S'(u)\nabla u = S'(u)\chi_{\{|u|\leq k\}}\nabla u = S'(u)T'_k(u)\nabla u = S'(u)\nabla T_k(u).$$

Derivation of the renormalized equation

Let u be a strong solution of the stochastic p-Laplace problem with \mathcal{F}_0 -measurable initial value $u_0 \in L^1(\Omega \times D)$, i.e.,

$$u(t)-u_0-\int_0^t \operatorname{div}\left(|\nabla u(s)|^{p-2}\nabla u(s)\right)ds=\int_0^t \Phi(s)\,d\beta_s.$$

For $S\in\mathcal{C}^2(\mathbb{R})$ with bounded derivatives, the Itô formula yields

$$S(u)(t) - S(u_0) - \int_0^t \operatorname{div} [S'(u)|\nabla u|^{p-2}\nabla u] \, ds$$

+ $\int_0^t S''(u)[|\nabla u|^p - \frac{1}{2}|\Phi|^2] \, ds = \int_D \Phi S'(u) \, d\beta_s$

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(R3) The energy dissipation condition holds true:

$$\lim_{k\to\infty}\mathbb{E}\int_{\{k<|u|< k+1\}}|\nabla u|^p\,dx\,dt=0.$$

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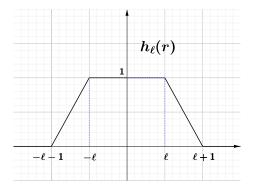
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Consistency of renormalized solutions

Let u be a renormalized solution with $\nabla u \in L^p(\Omega \times Q_T)^d$, $\ell > 0$.



Choosing

$$S_\ell(u) = \int_0^u h_\ell(r) \, dr$$

in the renormalized equation (R2), for $\ell \to \infty$, it follows that u is a strong solution.

L¹-contraction principle

Theorem

Let u, v be renormalized solutions with \mathcal{F}_0 -measurable initial values $u_0 \in L^1(\Omega \times D)$ and $v_0 \in L^1(\Omega \times D)$, respectively. Then,

$$\sup_{t\in[0,T]} \|u(t)-v(t)\|_{L^1(D)} \le \|u_0-v_0\|_{L^1(D)}$$

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a.s. in Ω .

Remark: As a consequence, renormalized solutions are unique.

Idea of the proof:

Let u, v be renormalized solutions with initial values u_0, v_0 . Subtracting the renormalized equation (*R*2) for v from the one for u we obtain

$$d[S(u) - S(v)] - \operatorname{div} [S'(u)|\nabla u|^{p-2}\nabla u - S'(v)|\nabla v|^{p-2}\nabla v] dt$$

+ $S''(u)|\nabla u|^p - S''(v)|\nabla v|^p dt - \frac{1}{2}\Phi^2(S''(u) - S''(v)) dt$
= $\Phi(S'(u) - S'(v)) d\beta_t.$

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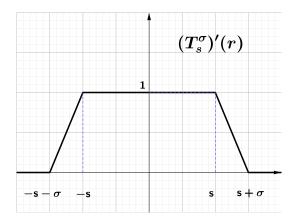
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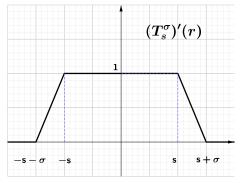
Formally, we choose S(r) = r and use the Itô formula for the absolute value.

Idea of the proof [Blanchard, Murat, Redwane; 2001]:

For $s, \sigma > 0$, we choose $S(r) = T_s^{\sigma}(r) := \int_0^r (T_s^{\sigma})'(\tau) d\tau$ with



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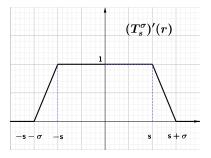


Itô formula for an approximation of the absolute value:

$$G_k(T_s^{\sigma}(u) - T_s^{\sigma}(v)) = \frac{1}{k} \int_0^{T_s^{\sigma}(u) - T_s^{\sigma}(v)} T_k(\tau) d\tau$$

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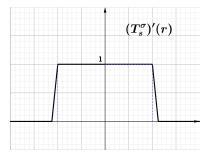
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Thanks to the energy dissipation condition (R3) we can pass to the limit with $\sigma \rightarrow 0$

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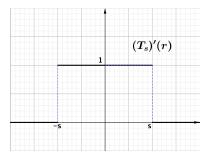
$$G_k(T_s^{\sigma}(u)-T_s^{\sigma}(v))=\frac{1}{k}\int_0^{T_s^{\sigma}(u)-T_s^{\sigma}(v)}T_k(\tau)\,d\tau.$$

Thanks to the energy dissipation condition (R3) we can pass to the limit with $\sigma \rightarrow 0$

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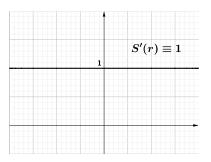
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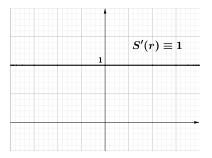


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Existence of renormalized solutions

Theorem (Existence)

For any \mathcal{F}_0 -measurable initial value $u_0 \in L^1(\Omega \times D)$, there exists a renormalized solution to the stochastic p-Laplace problem.



► Let $(u_0^n)_n \subset L^2(\Omega \times D)$ be an \mathcal{F}_0 -measurable sequence such that $u_0^n \to u_0$ in $L^1(\Omega \times D)$ for $n \to \infty$.

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- ► If u_n, u_m are strong solutions to the stochastic p-Laplace problem with initial values u₀ⁿ, u₀^m respectively, from the L¹-contraction principle we get

$$\sup_{t\in[0,T]} \|u_n(t) - u_m(t)\|_{L^1(D)} \le \|u_0^n - u_0^m\|_{L^1(D)}$$

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► There exists an *F_t*-adapted stochastic process *u* ∈ *L*¹(Ω; *C*([0, *T*]; *L*¹(*D*))) such that

$$\lim_{n\to\infty}u_n=u$$

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We claim that this function u is a renormalized solution with initial value u₀.

Passage to the limit in the renormalized equation

For all $S \in C^2(\mathbb{R})$ such that S' has compact support, the sequence of approximate solutions $(u_n) \subset L^1(\Omega \times Q_T)$, satisfies

$$S(u_n)(t) - S(u_0^n) - \int_0^t \operatorname{div} [S'(u_n) |\nabla u_n|^{p-2} \nabla u_n] \, ds$$

+ $\int_0^t S''(u_n) [|\nabla u_n|^p - \frac{1}{2} \Phi^2] \, ds = \int_0^t S'(u_n) \Phi \, d\beta_s$

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for all $t \in [0, T]$, a.s. in Ω . For the passage to the limit we need convergence of

 $\operatorname{div}\left[S'(u_n)|\nabla T_k(u_n)|^{p-2}\nabla T_k(u_n)\right] + S''(u_n)|\nabla T_k(u_n)|^p$

in $L^{p'}(0, t; W^{-1,p'}(D) + L^1(D))$, a.s. in Ω for $n \to \infty$ and any k > 0.

Strong convergence of the truncated gradients

 $\dot{\nabla}$...the method of [Blanchard; 1993] is to show that $(u_n)_{n\in\mathbb{N}}$ satisfies

$$\lim_{n,m\to\infty} \mathbb{E} \int_{Q_T} \left(|\nabla T_k(u_n)|^{p-2} \nabla T_k(u_n) - |\nabla T_k(u_m)|^{p-2} \nabla T_k(u_m) \right) \cdot \nabla (T_k(u_n) - T_k(u_m)) d(s,x) = 0.$$

As a consequence,

$$\lim_{n\to\infty}\nabla T_k(u_n)=\nabla T_k(u)$$

in $L^p(\Omega \times Q_T)^d$ for any k > 0, and we can pass to the limit with $n \to \infty$ in the renormalized equation.

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We need to find an equation for $Z(u_n - u_m)H(u_n)$ where H, Z are appropriate nonlinear test functions.

For $m, n \in \mathbb{N}$, u_0^n , $u_0^m \in L^2(\Omega \times D)$ \mathcal{F}_0 -measurable, $t \in [0, T]$

$$u_m(t) = u_0^m + \int_0^t \Delta_p(u_m) \, ds + \int_0^t \Phi \, d\beta_s$$
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$$\Longrightarrow (u_n - u_m)(t) = u_0^n - u_0^m + \int_0^t \Delta_p(u_n) - \Delta_p(u_m) ds$$

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$$\Longrightarrow (u_n-u_m)(t)=u_0^n-u_0^m+\int_0^t\Delta_p(u_n)-\Delta_p(u_m)\,ds.$$

Since $\ll u_n - u_m, u_n \gg_t = 0$, we obtain

$$(u_{n} - u_{m}, u_{n})_{2}(t) - (u_{0}^{n} - u_{0}^{m}, u_{0}^{n})_{2} + \int_{0}^{t} \langle \Delta_{p}(u_{n}) - \Delta_{p}(u_{m}), u_{n} \rangle \, ds$$
$$+ \int_{0}^{t} \langle \Delta_{p}(u_{n}), u_{n} - u_{m} \rangle \, ds = \int_{0}^{t} (\Phi, u_{n} - u_{m})_{2} \, d\beta_{s}$$

For $m, n \in \mathbb{N}$, $u_0^n, u_0^m \in L^2(\Omega \times D)$ \mathcal{F}_0 -measurable, let u_n, u_m be strong solutions to the *p*-Laplace evolution equation with initial value u_0^n, u_0^m respectively. For $Z \in W^{2,\infty}(\mathbb{R}), Z''$ piecewise continuous, $Z(0) = Z'(0) = 0, H \in C_b^2(\mathbb{R})$ it follows that

$$(Z(u_{n} - u_{m}), H(u_{n}))_{2}(t) = (Z(u_{0}^{n} - u_{0}^{m}), H(u_{0}^{n}))_{2}$$

+ $\int_{0}^{t} \langle \Delta_{p}(u_{n}) - \Delta_{p}(u_{m}), H(u_{n})Z'(u_{n} - u_{m}) \rangle ds$
+ $\int_{0}^{t} \langle \Delta_{p}(u_{n}), H'(u_{n})Z(u_{n} - u_{m}) \rangle ds$
+ $\frac{1}{2} \int_{0}^{t} \int_{D} \Phi^{2} H''(u_{n})Z(u_{n} - u_{m}) dx ds$
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$$(Z(u_n - u_m), H(u_n))_2(t) = (Z(u_0^n - u_0^m), H(u_0^n))_2 + \int_0^t \langle \Delta_p(u_n) - \Delta_p(u_m), H(u_n)Z'(u_n - u_m) \rangle \, ds + \int_0^t \langle \Delta_p(u_n), H'(u_n)Z(u_n - u_m) \rangle \, ds + \frac{1}{2} \int_0^t \int_D \Phi^2 H''(u_n)Z(u_n - u_m) \, dx \, ds + \int_0^t (\Phi H'(u_n), Z(u_n - u_m))_2 \, d\beta_s$$

► The unique renormalized solution u = u(t, r, u_r) of the stochastic p-Laplace equation starting in u_r at time r ∈ [0, T] is a time-homogeneous Markov process with transition semigroup

$$P_t(\varphi)(x) := \mathbb{E} \varphi(u(t,0,x)), \ t \in [0,T]$$

for bounded, measurable $\varphi : L^1(D) \to \mathbb{R}$.

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- $(P_t)_{t \in [0,T]}$ has the *Feller* property.
- (P_t)_{t∈[0,T]} has the *e*-property: For any φ : L¹(D) → ℝ bounded and Lipschitz continuous, x ∈ L¹(D) and ε > 0 there exists δ > 0 such that

$$|P_t(\varphi)(x) - P_t(\varphi)(z)| < \varepsilon$$

for all $z \in B(x, \delta)$ and all $0 \le t \le T$.

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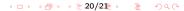
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► The distribution (P_x)_{x∈L¹(D)} of u(·, 0, x) is a Markov process on C([0, T]; L¹(D)).

Adding the first-order term

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with $F : \mathbb{R} \to \mathbb{R}^d$, convection effects enter into the problem.



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► General case:
$$F(u) \notin L^1(\Omega \times Q_T)$$
, however $F(T_k(u)) \in L^1(\Omega \times Q_T)$ for all $k > 0$.



